



## A multi-messenger probe of the nature of neutrino mass

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in collaboration with
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arXiv: 2212.00737 [astro-ph.HE]

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Texas A&M University, College Station

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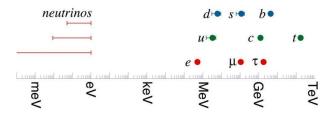
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 $\Longrightarrow \textbf{Nonzero Neutrino Mass} \Longrightarrow \textbf{BSM Physics}$ 

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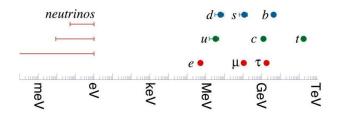


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#### $\Longrightarrow$ Nonzero Neutrino Mass $\Longrightarrow$ BSM Physics



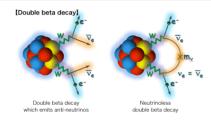
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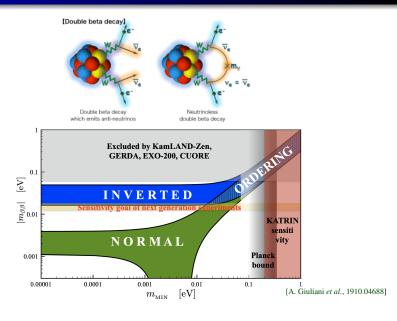
 ${\bf Majorana\ or\ Dirac\ (or\ something\ in\ between)?}$ 

Only experiments can tell.

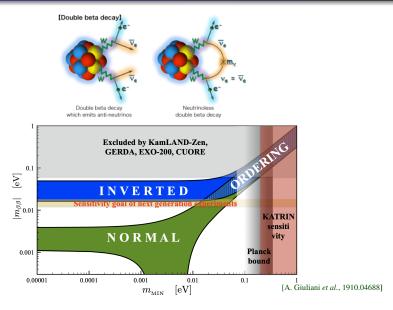
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What if the Majorana mass is small?

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- If  $M_R = 0$ , lepton number is preserved and neutrinos are **Dirac**.
- If  $M_R \neq 0$ , neutrinos are **Majorana**.
- If  $||M_R|| \ll ||m_D||$ , neutrinos are **pseudo-Dirac** (small active-sterile mass splitting).
- But isn't it more natural to have  $||M_R|| \gg ||m_D||$  (seesaw)? [Minkowski (PLB '77); Mohapatra, Senjanovic (PRL '80); Yanagida '79; Gell-Mann, Ramond, Slansky '79]

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- Maybe, but  $||M_R|| \ll ||m_D||$  is a logical possibility too. [Wolfenstein (NPB '81); Petcov (PLB '82); Valle (PRD '82); Valle, Singer (PRD '83); Kobayashi, Lim (PRD '01)]

- A good starting point: Dirac neutrino models with  $m_D$  naturally small and  $M_R=0$  at renormalizable level, e.g. in **Dirac seesaw**. [Roncadelli, Wyler (PLB '83); Roy, Shanker (PRL '84); Dick, Lindner, Ratz, Wright (PRL '00); Murayama, Pierce (PRL '02); Gu, He (JCAP '06); Joshipura, Mohanty, Pakvasa (PRD '14); Ma, Srivastava (PLB '15); Ma, Popov (PLB '17); Earl, Fong, Gregoire, Tonero (JCAP '20); ...]
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- Weinberg operators  $(\Psi \Psi H H)/M_{\rm Pl}$  and  $(\Psi' \Psi' H' H')/M_{\rm Pl}$  induce small diagonal  $M_{\nu}$  entries  $\Longrightarrow$  **Pseudo-Dirac neutrinos**.
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- $\bullet$  But excluded by BBN and solar neutrino constraints (for maximal mixing).  $\ensuremath{\boxdot}$

$$\delta m^2 \lesssim 10^{-8}~\text{eV}^2~\text{from BBN, [Barbieri, Dolgov (PLB '90)]} \\ 10^{-11}~\text{eV}^2~\text{from solar. [de Gouvêa, Huang, Jenkins (PRD '09); Ansarifard, Farzan (PRD '23)]}$$

- An alternative is to gauge B L.
- Introduce a singlet scalar S carrying two units of B-L.
- Lowest-order quantum gravity corrections are of the form  $(\Psi \Psi HHS)/M_{\rm Pl}^2$ .
- For  $\langle S \rangle = v_{BL}$ , leads to diagonal elements of  $M_{\nu}$  of order  $v^2 v_{BL}/M_{\rm Pl}^2$ .
- For  $v_{BL} = (10^4 10^{14}) \text{ GeV}$ , generates  $\delta m^2 \sim (10^{-22} 10^{-12}) \text{ eV}^2$ .
- Consistent with solar neutrino data. ©

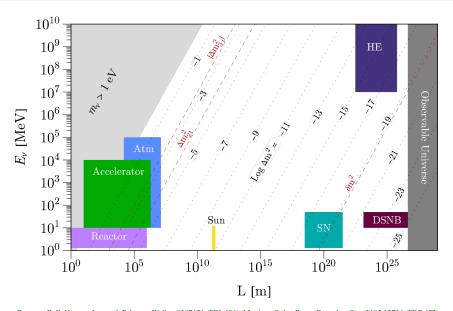
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### How to probe these tiny $\delta m^2$ values?

Oscillation effects are suppressed, unless L and E are such that  $\delta m^2 L/E \sim 1$ .

# Need astrophysical baselines



Beacom, Bell, Hooper, Learned, Pakvasa, Weiler, 0307151 (PRL '04); Martinez-Soler, Perez-Gonzalez, Sen, 2105.12736 (PRD '22)

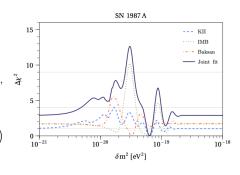
$$\begin{split} P_{aa}(E_{\nu}) &= \frac{1}{2} \left( 1 + e^{-\left(\frac{L}{L_{\rm coh}}\right)^2} \cos\left(\frac{2\pi L}{L_{\rm osc}}\right) \right). \\ L_{\rm osc} &= \frac{4\pi E_{\nu}}{\delta m^2} \approx 20 \ \mathrm{kpc} \left(\frac{E_{\nu}}{25 \ \mathrm{MeV}}\right) \left(\frac{10^{-19} \ \mathrm{eV}^2}{\delta m^2}\right), \\ L_{\rm coh} &= \frac{4\sqrt{2} E_{\nu}}{|\delta m^2|} (E_{\nu} \sigma_x) \\ &\approx 114 \ \mathrm{kpc} \left(\frac{E_{\nu}}{25 \ \mathrm{MeV}}\right)^2 \left(\frac{10^{-19} \ \mathrm{eV}^2}{\delta m^2}\right) \left(\frac{\sigma_x}{10^{-13} \ \mathrm{m}}\right) \end{split}$$

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Martinez-Soler, Perez-Gonzalez, Sen, 2105.12736 (PRD '22)

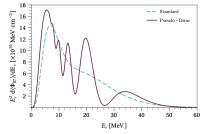
 $E_v$  [MeV]

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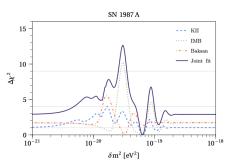


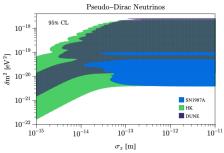
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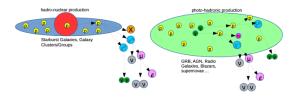


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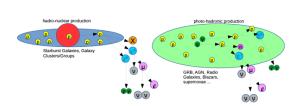


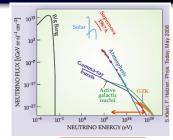
# Astrophysical neutrinos



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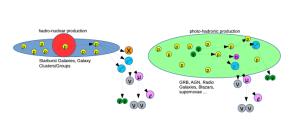
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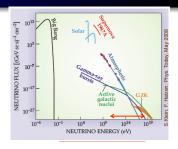




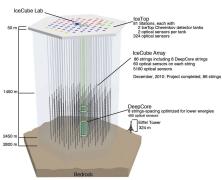
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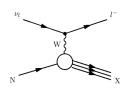
### Need gigantic detectors to compensate for the tiny flux.



(

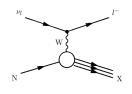
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$$(\nu_e : \nu_\mu : \nu_\tau) = (1:2:0)_{\bigstar} \longrightarrow (1:1:1)_{\oplus}.$$



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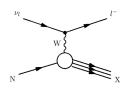




CC EM/NC all (shower)

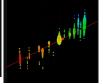
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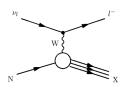


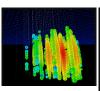


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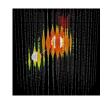




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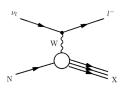


CC tau 'double bang' (only at  $E_{\nu} \gtrsim 100~{\rm TeV})$ 

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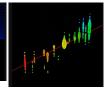
Flavor composition:

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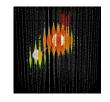




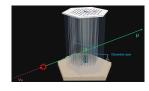
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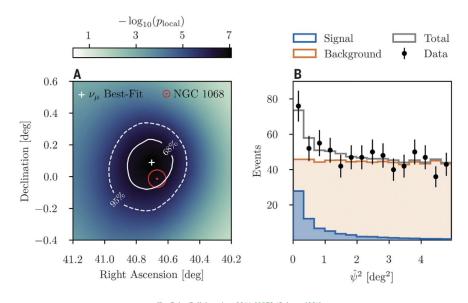


Throughgoing muon (track only, huge statistics)

[Picture courtesy: C. Kopper]

**Showers:** Good energy resolution, but poor angular resolution **Tracks:** Excellent angular resolution ( $< 1^{\circ}$ ), but modest energy resolution (< 30%)

Track events are ideal for astrophysical source identification.

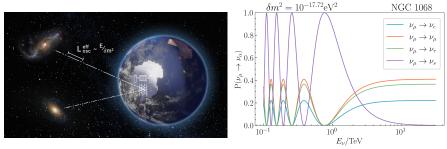


[IceCube Collaboration, 2211.09972 (Science '22)]

# A new probe of pseudo-Dirac neutrinos



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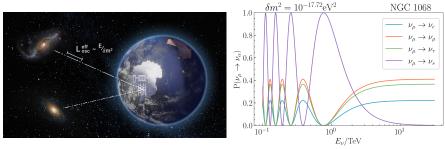


Carloni, Martínez-Soler, Argüelles, Babu, BD, 2212.00737

Oscillation probability:

$$\begin{split} P_{\alpha\beta} &= \frac{1}{2} \sum_{j=1}^{3} |U_{\beta j}|^2 |U_{\alpha j}|^2 \left[ 1 + \cos\left(\frac{\delta m_j^2 L_{\rm eff}}{2E_\nu}\right) \right], \\ \text{with } L_{\rm eff} &= \int \frac{dz}{H(z)(1+z)^2} \text{ and } H(z) = H_0 \sqrt{\Omega_m (1+z)^3 + \Omega_\Lambda + (1-\Omega_m - \Omega_\Lambda)(1+z)^2}. \end{split}$$

## A new probe of pseudo-Dirac neutrinos



Carloni, Martínez-Soler, Argüelles, Babu, BD, 2212.00737

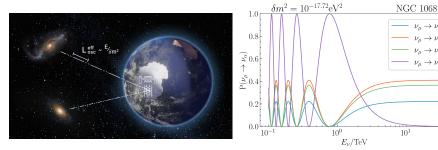
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with 
$$L_{\rm eff}=\int \frac{dz}{H(z)(1+z)^2}$$
 and  $H(z)=H_0\sqrt{\Omega_m(1+z)^3+\Omega_\Lambda+(1-\Omega_m-\Omega_\Lambda)(1+z)^2}.$ 

 $\bullet \ \ \text{Typical oscillation length:} \ L_{\rm osc} = \tfrac{2E_{\nu}}{\delta m^2} \approx 6.4 \ {\rm Mpc} \left( \tfrac{E_{\nu}}{1 \ {\rm TeV}} \right) \left( \tfrac{2 \times 10^{-18} \ {\rm eV}^2}{\delta m^2} \right) .$ 

## A new probe of pseudo-Dirac neutrinos



Carloni, Martínez-Soler, Argüelles, Babu, BD, 2212.00737

Oscillation probability:

$$P_{\alpha\beta} = \frac{1}{2} \sum_{j=1}^{3} |U_{\beta j}|^2 |U_{\alpha j}|^2 \left[ 1 + \cos\left(\frac{\delta m_j^2 L_{\text{eff}}}{2E_{\nu}}\right) \right],$$

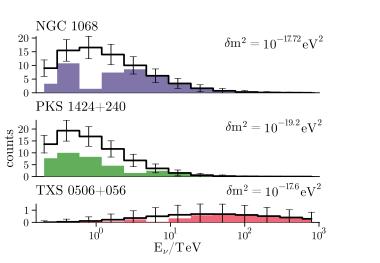
with 
$$L_{\mathrm{eff}}=\int \frac{dz}{H(z)(1+z)^2}$$
 and  $H(z)=H_0\sqrt{\Omega_m(1+z)^3+\Omega_\Lambda+(1-\Omega_m-\Omega_\Lambda)(1+z)^2}.$ 

- Typical oscillation length:  $L_{\rm osc} = \frac{2E_{\nu}}{\delta m^2} \approx 6.4 \ {\rm Mpc} \left(\frac{E_{\nu}}{1 \ {\rm TeV}}\right) \left(\frac{2 \times 10^{-18} \ {\rm eV}^2}{\delta m^2}\right)$ .
- Typical coherence length: [Kersten, Smirnov, 1512.09068 (EPJC '16); Rink, Sen, 2211.16520]  $L_{\rm coh} = \frac{4\sqrt{2}E_{\nu}^2}{|\delta m^2|} \approx 10^{10} \,{\rm Mpc} \left(\frac{E_{\nu}}{1 \,{\rm TeV}}\right)^2 \left(\frac{2\times 10^{-18} \,{\rm eV}^2}{|\delta m^2|}\right) \left(\frac{\sigma_x}{10^{-10} \,{\rm m}}\right) \gg L_{\rm osc}.$

 $10^{1}$ 

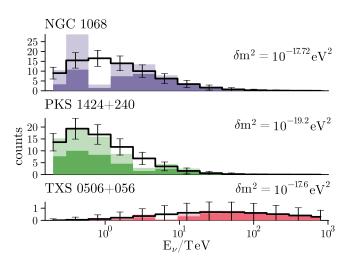
### **Event Distributions**

Source	Source Type	$-\log_{10} p_{\mathrm{local}}$	$\hat{n}_s$	$\hat{\gamma}$	z
NGC 1068	SBG/AGN	$7.0 (5.2\sigma)$	79	3.2	0.0038 (16 Mpc)
PKS 1424+240	BLL	$4.0 \ (3.7\sigma)$	77	3.5	0.6047 (2.6 Gpc)
TXS 0506+056	BLL/FSRQ	$3.6 (3.5\sigma)$	5	2.0	0.3365 (1.4 Gpc)

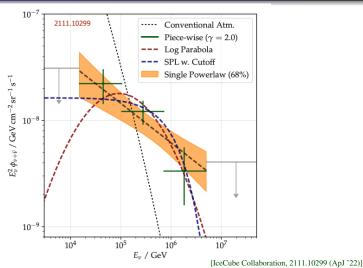


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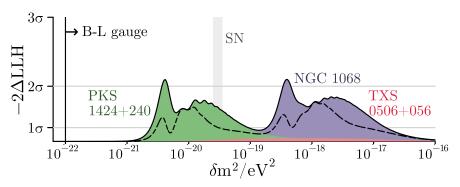
### Power-law flux



$$\Phi(E_{\nu}) = \begin{cases}
\Phi_0(E_{\nu}/E_0)^{-\gamma} & \text{(single power law)} \\
\Phi'_0(E_{\nu}/E_0)^{-\gamma'} e^{-E_{\nu}/E_{\text{cutoff}}} & \text{(SPL with cutoff)} \\
\Phi''_0 \cdot (E/E_0)^{-(\alpha+\beta\log(E/E_0))} & \text{(log parabola)}
\end{cases}$$

### First IceCube constraints on $\delta m^2$

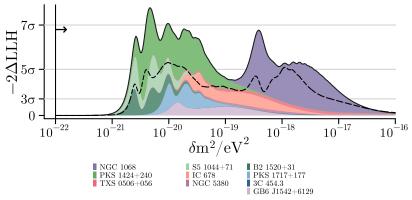
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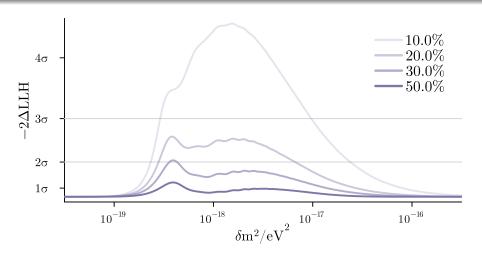
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## Future IceCube-Gen2 sensitivity

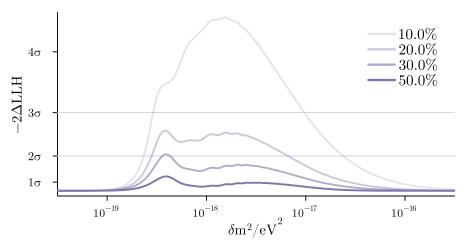
Source	Source Type	$-\log_{10} p_{local}$	$\hat{n}_{S}$	Ŷ	z
NGC 1068	SBG/AGN	7.0	79	3.2	0.0038
PKS 1424+240	BLL	4.0	77	3.5	0.6047
TXS 0506+056	BLL/FSRQ	3.6	5	2.0	0.3365
S5 1044+71	FSRQ	1.3	45	4.3	1.1500
IC 678	GAL	0.9	22	3.1	0.04799
NGC 5380	GAL	0.9	4	2.4	0.010584
B2 1520+31	FSRQ	1.0	35	4.3	1.48875
PKS 1717+177	BLL	1.0	34	4.3	0.137
3C 454.3	FSRQ	1.2	1	1.5	0.859
GB6 J1542+6129	BLL	1.9	16	4.3	0.117



# Effect of energy resolution



# Effect of energy resolution



- Lot of room for improvement with better technology.
- Showers can do much better, if the angular resolution can be improved somehow.
- Possible with KM3NeT or P-ONE (or something even bigger and better).



#### Conclusion

- The nature of neutrino mass (Majorana, Dirac, or pseudo-Dirac) has to be experimentally determined.
- We proposed a new experimental probe of pseudo-Diracness of neutrinos using high-energy astrophysical neutrinos.
- Made possible by recent breakthroughs in multi-messenger neutrino astrophysics.
- Current IceCube data on the three most significant astrophysical neutrino sources already constrain  $\delta m^2$  in the range of  $10^{-21}-10^{-16}~{\rm eV^2}$  with up to  $2\sigma$  significance.
- With additional sources and more statistics at IceCube-Gen2 (+KM3NeT), a larger range of  $\delta m^2$  can be probed with higher significance.
- Robust against astrophysical flux and flavor ratio uncertainties.

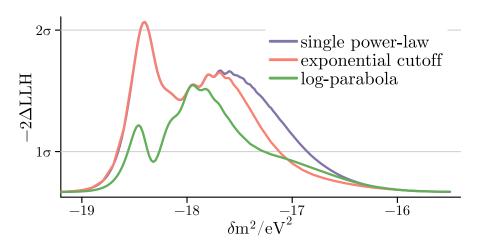
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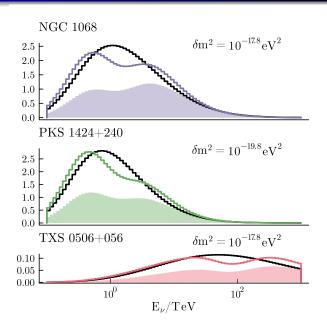
### Thank You.



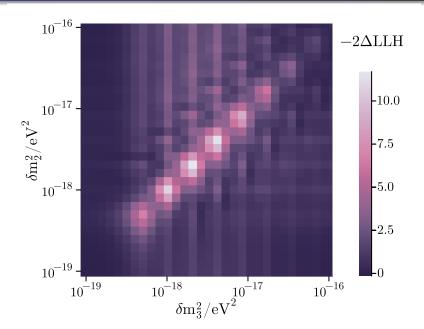
### Different flux models



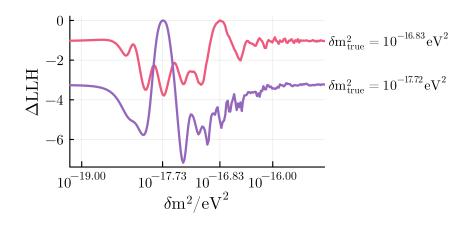
## Different binning



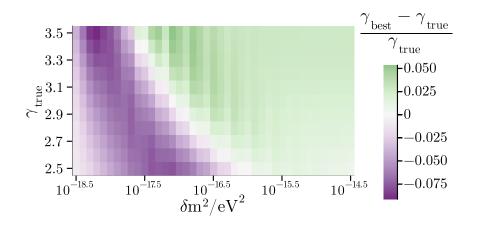
# Different mass splittings



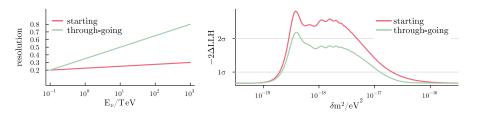
# Recovery of true pseudo-Dirac parameters



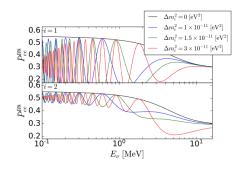
# Standard Model fits to pseudo-Dirac reality



# Linear models for energy resolution



### Solar neutrino constraint



### Solar neutrino constraint

