

Comprehensive treatment of correlations at different energy scales in nuclei using Green's functions

Lecture 1: 8/28/07	Propagator description of single-particle motion and the link with experimental data
Lecture 2: 8/29/07	From Hartree-Fock to spectroscopic factors < 1 : inclusion of long-range correlations
Lecture 3: 8/29/07	Role of short-range and tensor correlations associated with realistic interactions
Lecture 4: 8/30/07	Dispersive optical model and predictions for nuclei towards the dripline
Adv. Lecture 1: 8/30/07	Saturation problem of nuclear matter & pairing in nuclear and neutron matter
Adv. Lecture 2: 8/31/07	Quasi-particle density functional theory

Wim Dickhoff
Washington University in St. Louis

Outline

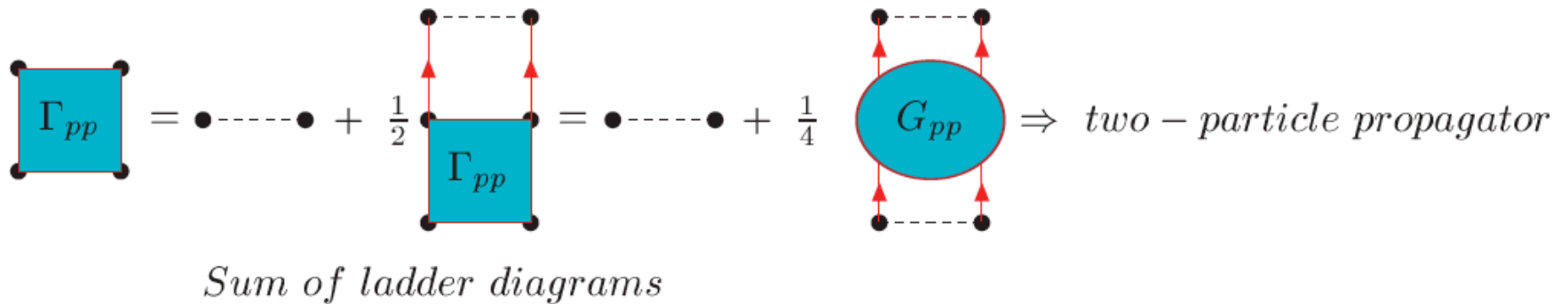
- SRC for free particles
- Ladders in the medium
- Self-energy and Dyson equation
- Nuclear matter simplifications
- Self-consistent Green's functions in nuclear matter & results
- SRC in finite nuclei: where are the high-momentum nucleons
- Summary of sp strength in closed-shell nuclei
- Other nuclei
- N very different from Z
- Nuclear matter with isospin polarization

Short-range correlations for two free particles

Solve the Schrödinger equation or the equivalent "T"-matrix

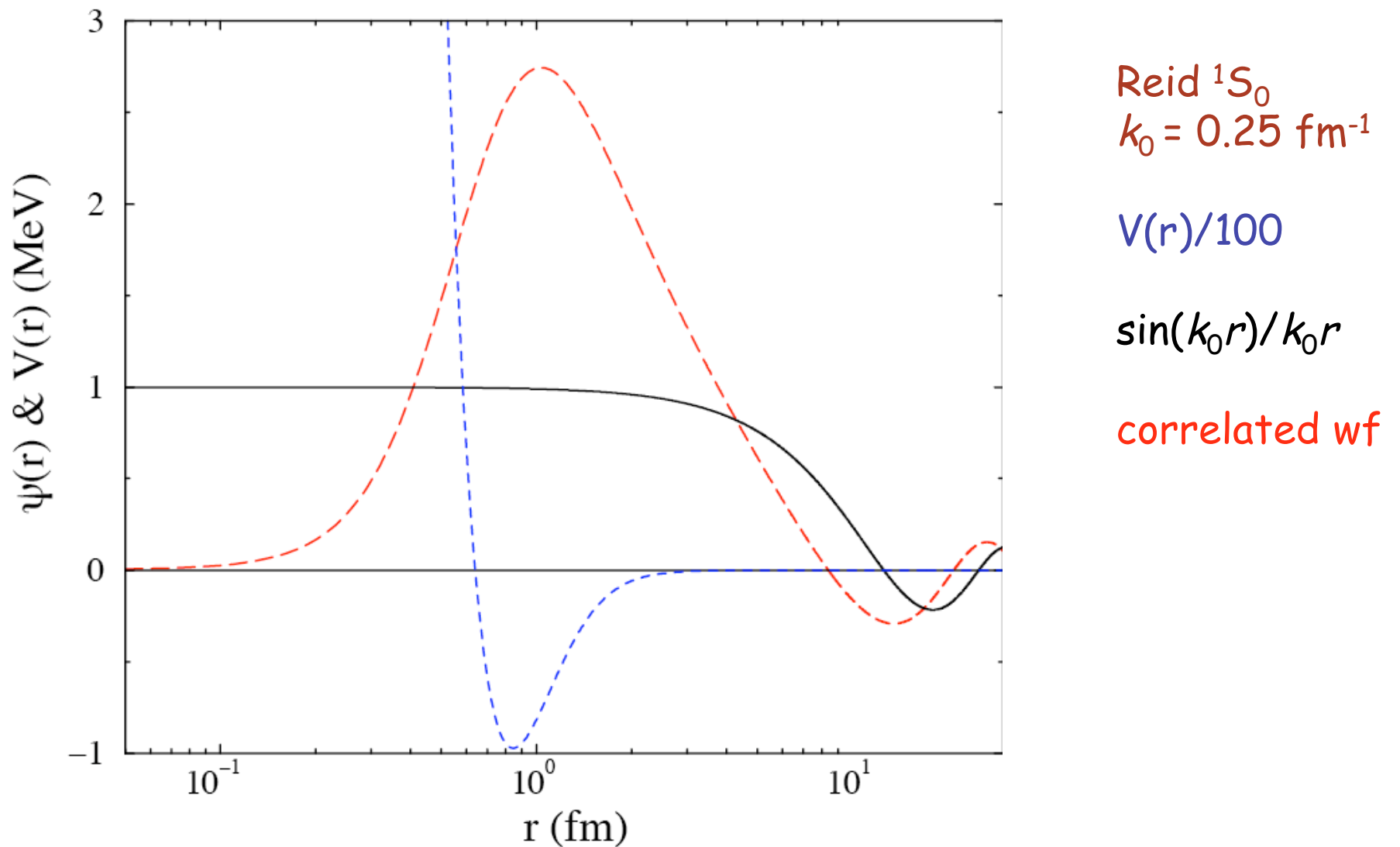
$$\langle k\ell | \Gamma_{pp}^{JST}(k_0) | k'\ell' \rangle = \langle k\ell | V^{JST} | k'\ell' \rangle + \frac{m}{2\hbar^2} \sum_{\ell''} \int_0^\infty \frac{dq}{(2\pi)^3} q^2 \langle k\ell | V^{JST} | q\ell'' \rangle \frac{1}{k_0^2 - q^2 + i\eta} \langle q\ell'' | \Gamma_{pp}^{JST}(k_0) | k'\ell' \rangle$$

Effective interaction



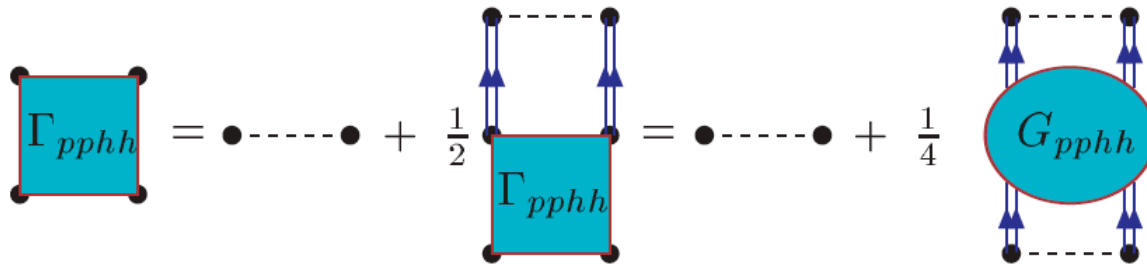
Sum of ladder diagram takes care of SRC
Also in the medium!

Relative wave function and potential



Ladder diagrams in the medium (options)

Ladders in the medium



\Rightarrow self-energy calculation

$$\langle k\ell | \Gamma_{pphh}^{JST}(K, E) | k' \ell' \rangle = \langle k\ell | V^{JST} | k' \ell' \rangle + \frac{1}{2} \sum_{\ell''} \int_0^{\infty} \frac{dq}{(2\pi)^3} q^2 \langle k\ell | V^{JST} | q\ell'' \rangle G_{pphh}^f(q; K, E) \langle q\ell'' | \Gamma_{pphh}^{JST}(K, E) | k' \ell' \rangle$$

G_{pphh}^f has different form depending on the level of sophistication

Nuclear matter:

$$G_{BG}^f(k_1, k_2; E) = \frac{\theta(k_1 - k_F) \theta(k_2 - k_F)}{E - \varepsilon(k_1) - \varepsilon(k_2) + i\eta}$$

Bethe-Goldstone

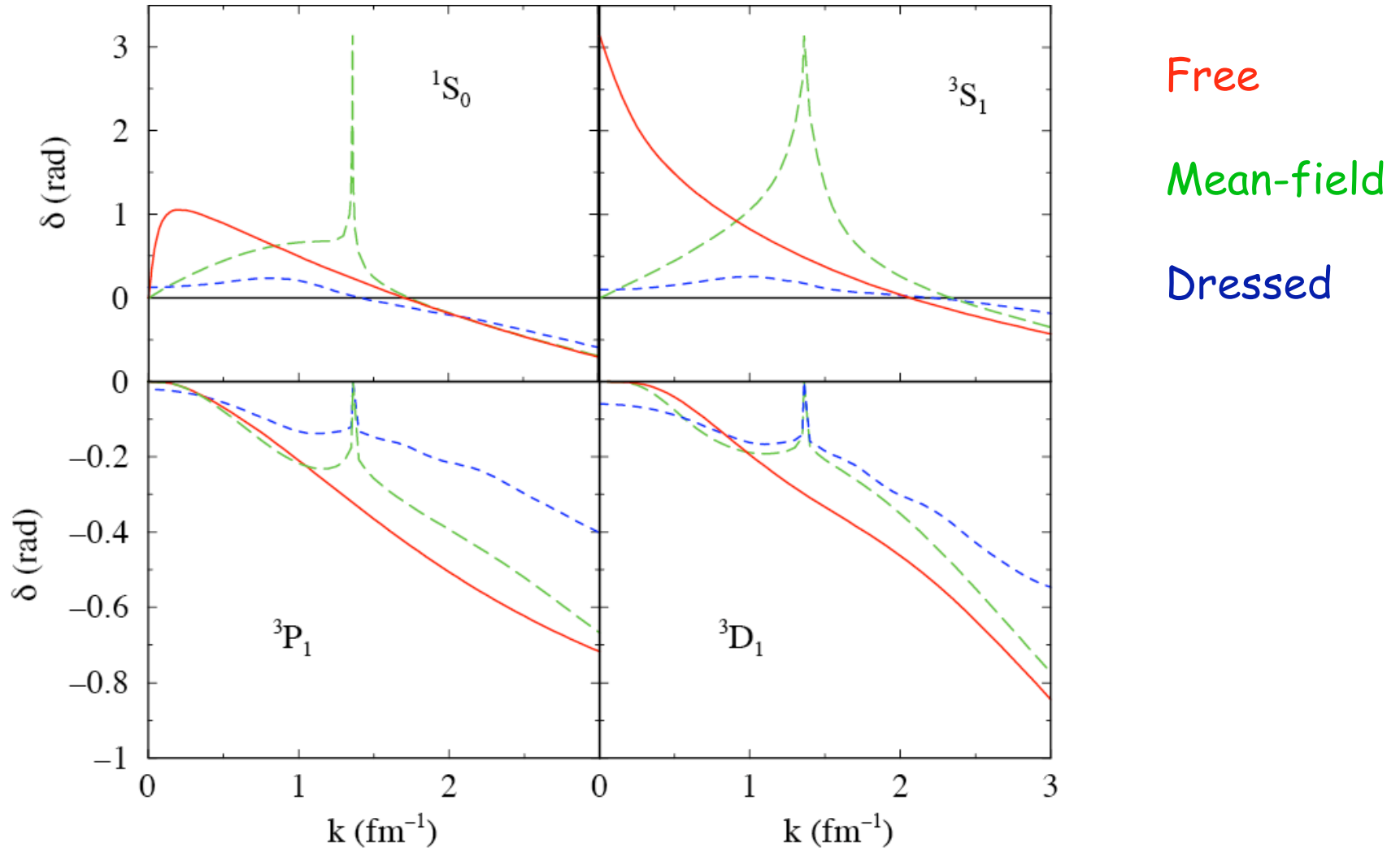
$$G_{GF}^f(k_1, k_2; E) = \frac{\theta(k_1 - k_F) \theta(k_2 - k_F)}{E - \varepsilon(k_1) - \varepsilon(k_2) + i\eta} - \frac{\theta(k_F - k_1) \theta(k_F - k_2)}{E - \varepsilon(k_1) - \varepsilon(k_2) - i\eta}$$

Galitskii-Feynman

$$G_{pphh}^f(k_1, k_2; E) = \int_{\varepsilon_F}^{\infty} dE_1 \int_{\varepsilon_F}^{\infty} dE_2 \frac{S_p(k_1; E_1) S_p(k_2; E_2)}{E - E_1 - E_2 + i\eta} - \int_{-\infty}^{\varepsilon_F} dE_1 \int_{-\infty}^{\varepsilon_F} dE_2 \frac{S_h(k_1; E_1) S_h(k_2; E_2)}{E - E_1 - E_2 - i\eta} \quad \text{SCGF}$$

Green's functions III 5

Phase shifts for dressed nucleons



PRC60, 064319 (1999) also PRC58, 2807 (1998) *Green's functions III 6*

Dyson equation and spectral functions in nuclear matter (some simplifications)

$$G(k; E) = G^{(0)}(k; E) + G^{(0)}(k; E)\Sigma(k; E)G(k; E)$$

Dyson equation

$$= \frac{1}{E - \varepsilon(k) - \Sigma(k; E)}$$

$$G^{(0)}(k; E) = \frac{\theta(k - k_F)}{E - \varepsilon(k) + i\eta} + \frac{\theta(k_F - k)}{E - \varepsilon(k) - i\eta}$$

Noninteracting sp propagator

$$S_p(k; E) = -\frac{1}{\pi} \frac{\text{Im}\Sigma(k; E)}{(E - \varepsilon(k) - \text{Re}\Sigma(k; E))^2 + (\text{Im}\Sigma(k; E))^2}$$

particle
spectral function

$$S_h(k; E) = \frac{1}{\pi} \frac{\text{Im}\Sigma(k; E)}{(E - \varepsilon(k) - \text{Re}\Sigma(k; E))^2 + (\text{Im}\Sigma(k; E))^2}$$

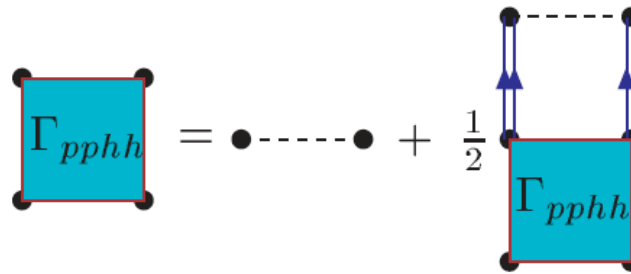
hole
spectral function

$$G(k; E) = \int_{\varepsilon_F}^{\infty} dE' \frac{S_p(k; E')}{E - E' + i\eta} + \int_{-\infty}^{\varepsilon_F} dE' \frac{S_h(k; E')}{E - E' - i\eta}$$

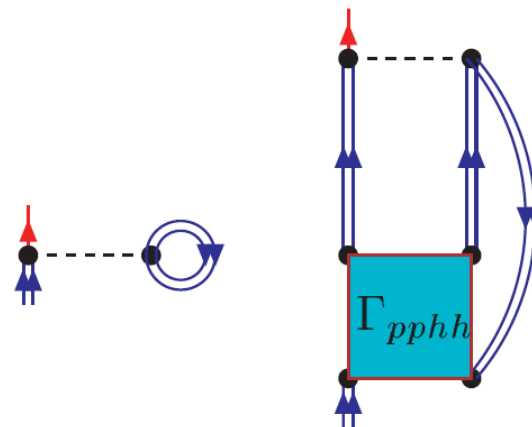
numerator sp strength
denominator where

Green's functions III 7

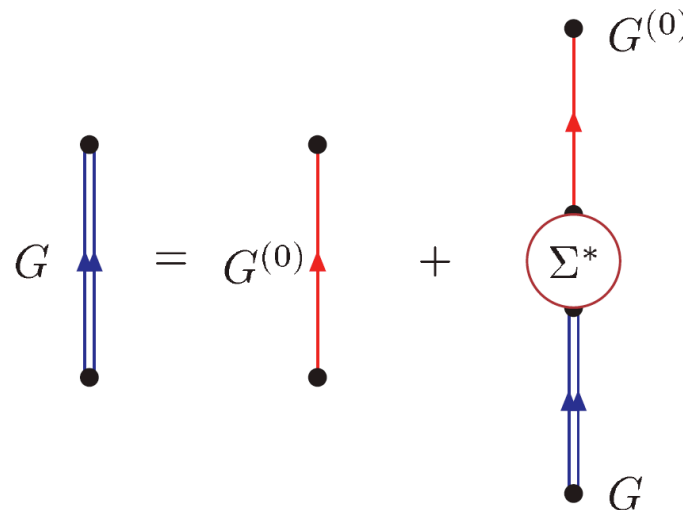
Self-consistency



Interaction



Self-energy



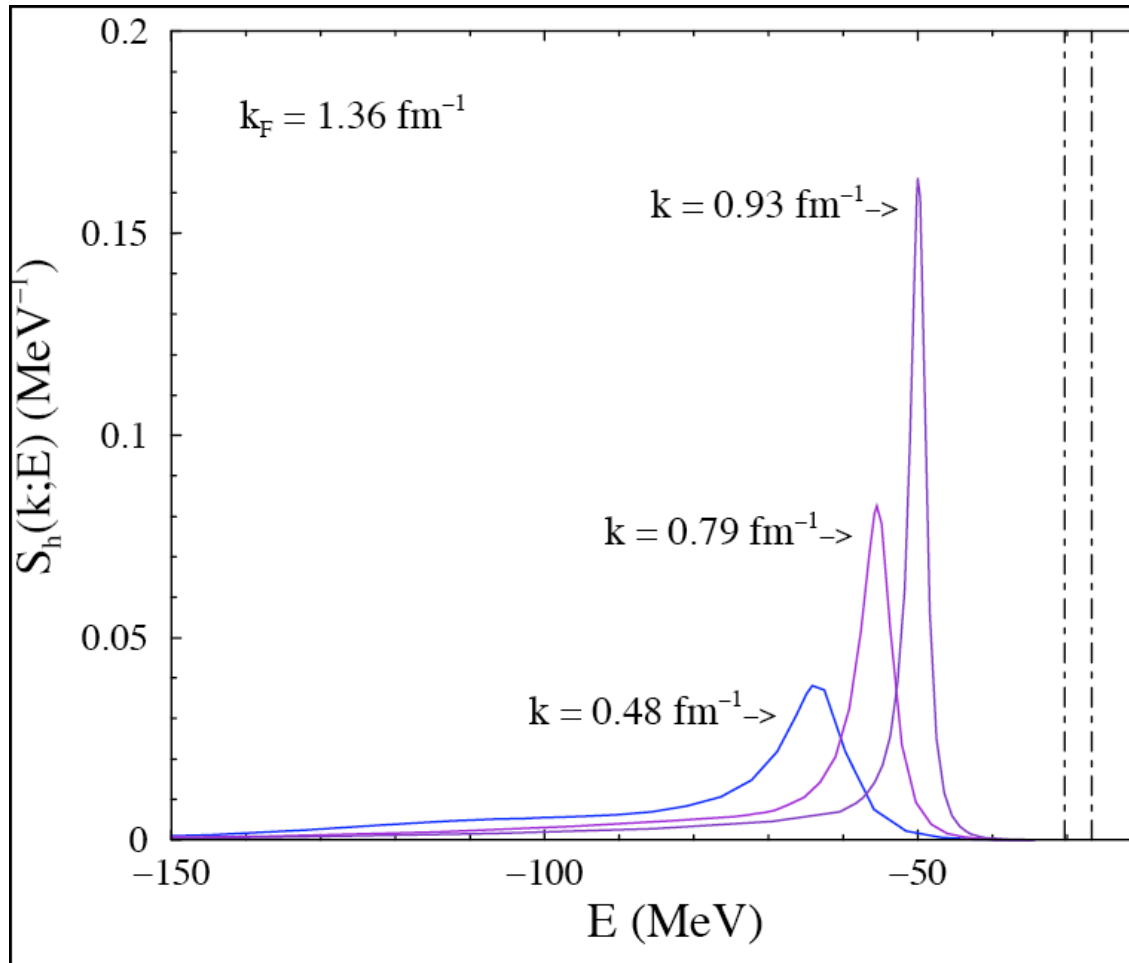
Dyson equation

Green's functions III 8

Groups

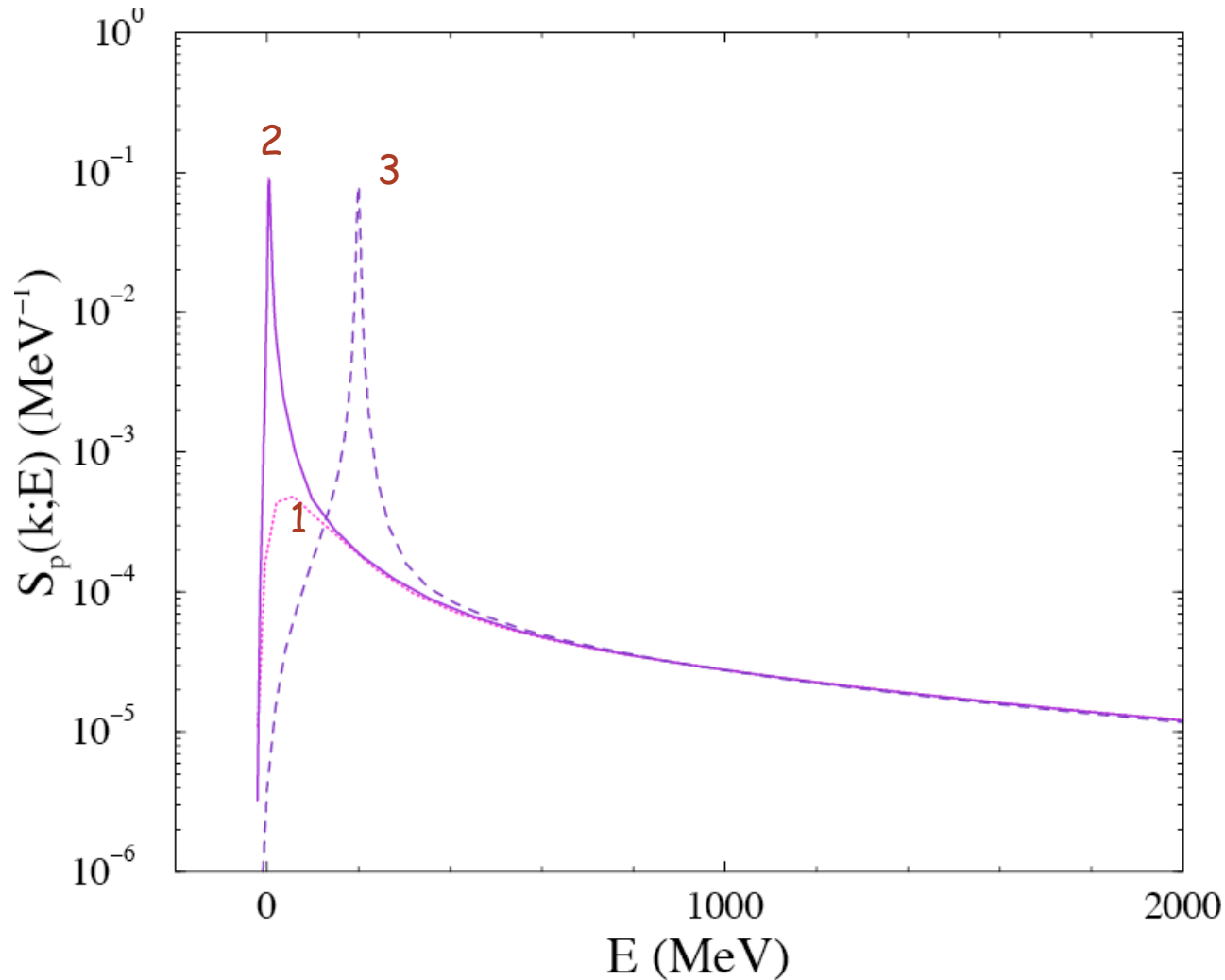
- St. Louis ⇒ Complete self-consistency for spectral functions
(Ramos, Vonderfecht, Gearhart, Roth/Stoddard)
- Ghent ⇒ Discrete method (Van Neck, Dewulf)
- Cracow ⇒ Separable & soft interactions (Bozek, Czerski, Soma)
- Tübingen ⇒ Finite temperature & soft interactions (Müther, Frick)
- Barcelona ⇒ Finite temperature & soft interactions (Polls, Ramos, Rios)
- Giessen ⇒ Interaction related to cross section (Lenske et al)

Illustrative results for mean-field input



Strength distribution as in nuclei ...
peak strength +10% at lower energy + global depletion!!

Where does the strength go?



All tails the same!
 \Rightarrow SRC

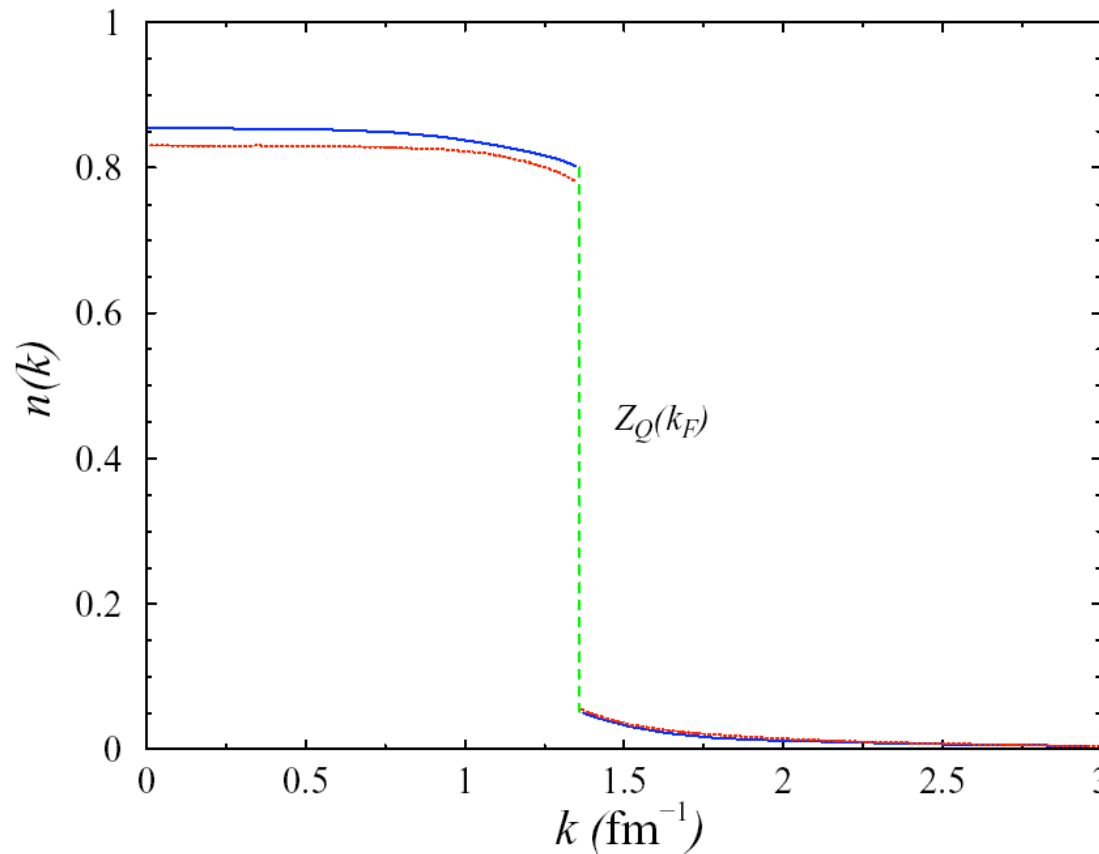
$$k_1 = 0.79 \text{ fm}^{-1}$$
$$k_2 = 1.74 \text{ fm}^{-1}$$
$$k_3 = 3.51 \text{ fm}^{-1}$$

$k < k_F$: 17% $> \varepsilon_F$ with 13% above 100 MeV (7% above 500 MeV)

Without tensor force only 10.5% above ε_F Green's functions III 11

Short-range correlations in nuclear matter and $n(k)$

$n(k=0) = 0.83 / 0.85 \Rightarrow$ finite nuclei



$$n(k) = \int_{-\infty}^{\varepsilon_F} dE S_h(k; E)$$

Reid soft core
 $k_F = 1.36 \text{ fm}^{-1}$

Old prediction!
 New result

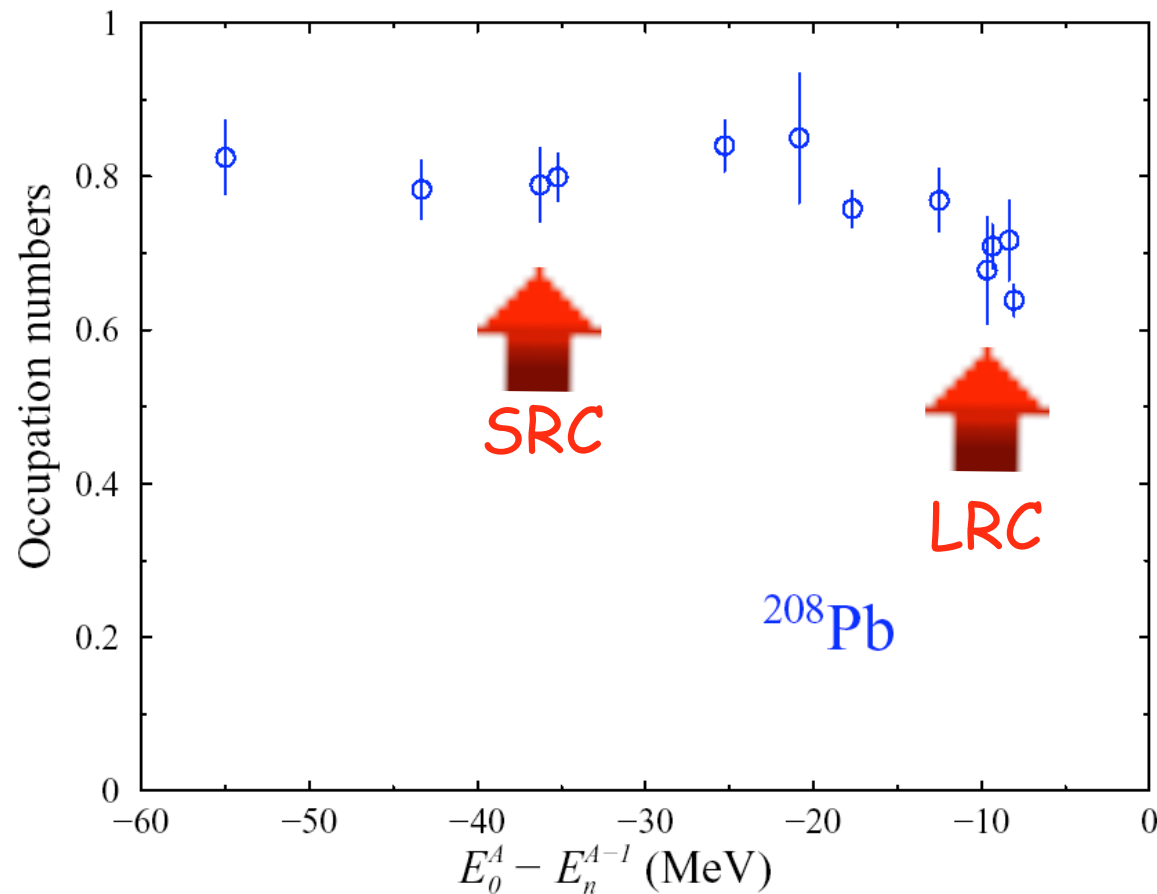
$Z_Q(k_F) = 0.72$

$Z_Q(k_F) = 0.75$

B.E.Vonderfecht et al. Nucl. Phys. A555, 1 (1993)
 E.R.Stoddard, thesis (self-consistent ladders)

M. van Batenburg (thesis, 2001) & L. Lapikás from $^{208}\text{Pb} (e, e' p) ^{207}\text{Tl}$

Occupation of deeply-bound proton levels from EXPERIMENT



Up to 100 MeV
missing energy
and
270 MeV/c
missing momentum

Covers the whole
mean-field domain
for the FIRST time!!

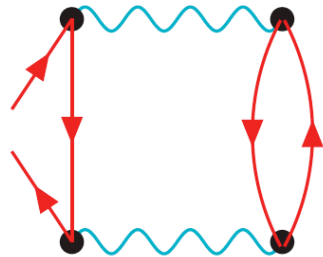
Confirmation of theory

Two effects associated with short-range correlations

- Depletion of the Fermi sea
- Admixture of high-momentum components to replace depleted strength

Location of high-momentum components

high momenta



require specific intermediate states

External line k (large).

Intermediate holes $< k_F$, say total momentum ~ 0 .

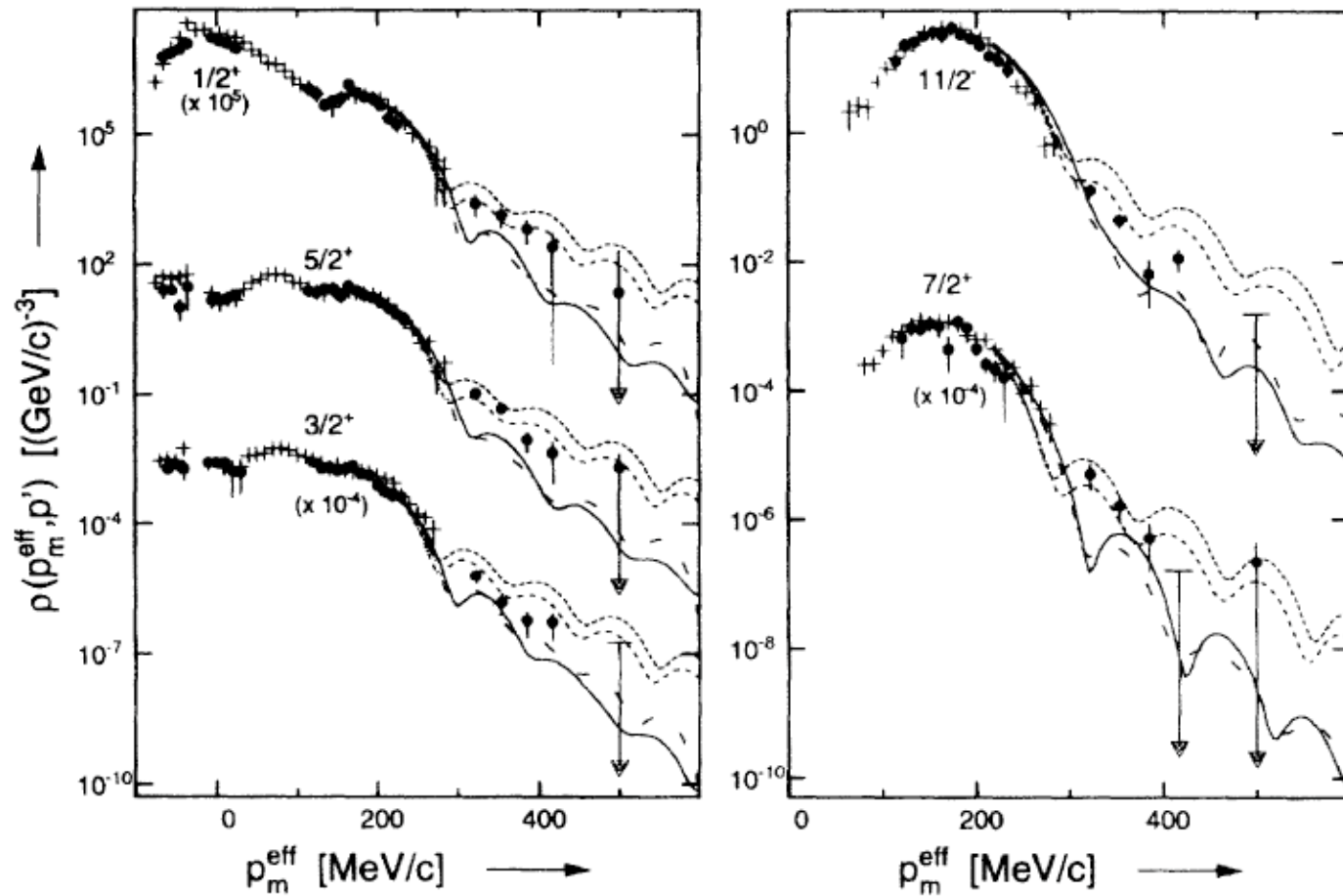
Momentum conservation: intermediate particle $-k$

\Rightarrow Energy intermediate state $\sim \langle \varepsilon_{2h} \rangle - \varepsilon(k)$

\Rightarrow the higher k the more negative the location of its strength

\Rightarrow no high-momentum components near ε_F

High-momenta near ε_F ?

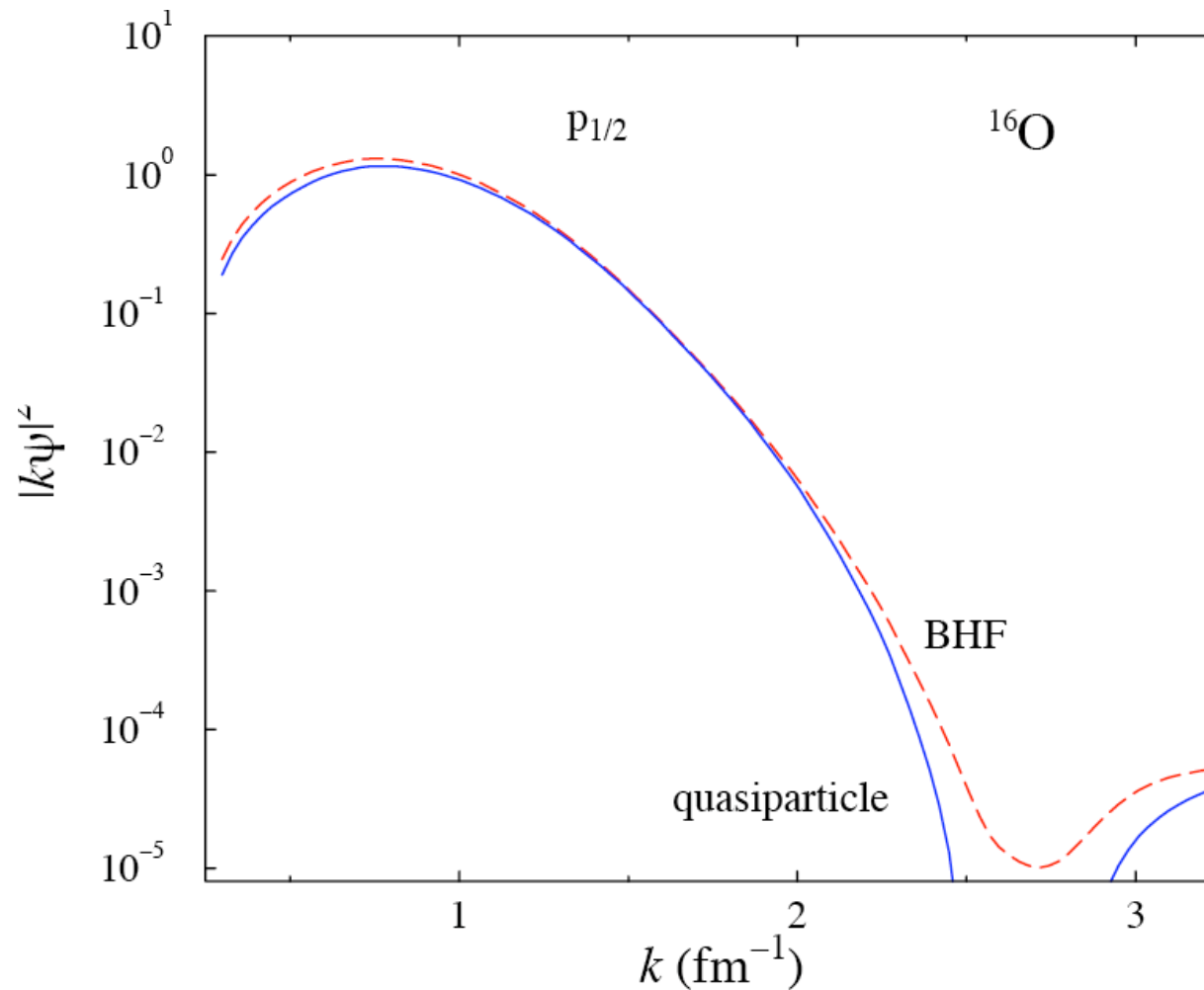


NO!

I. Bobeldijk *et al.*, *Phys. Rev. Lett.* **73**, 2684 (1994)

Green's functions
III16

SRC (only) calculated in ^{16}O



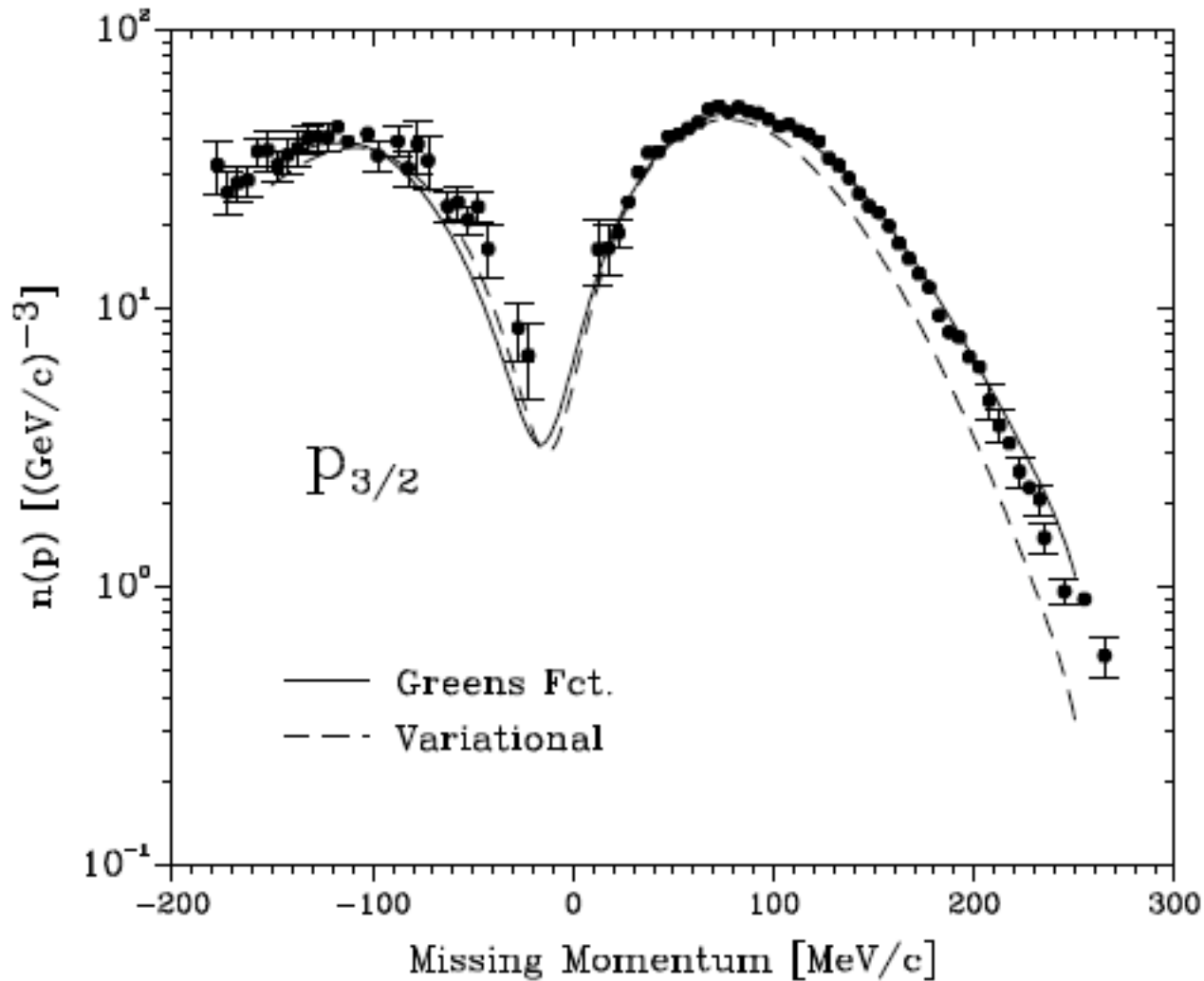
No enhancement
of high k near ε_F

Not observed
experimentally
either!

PRL73,2684(1994)
PLB344,85(1995)

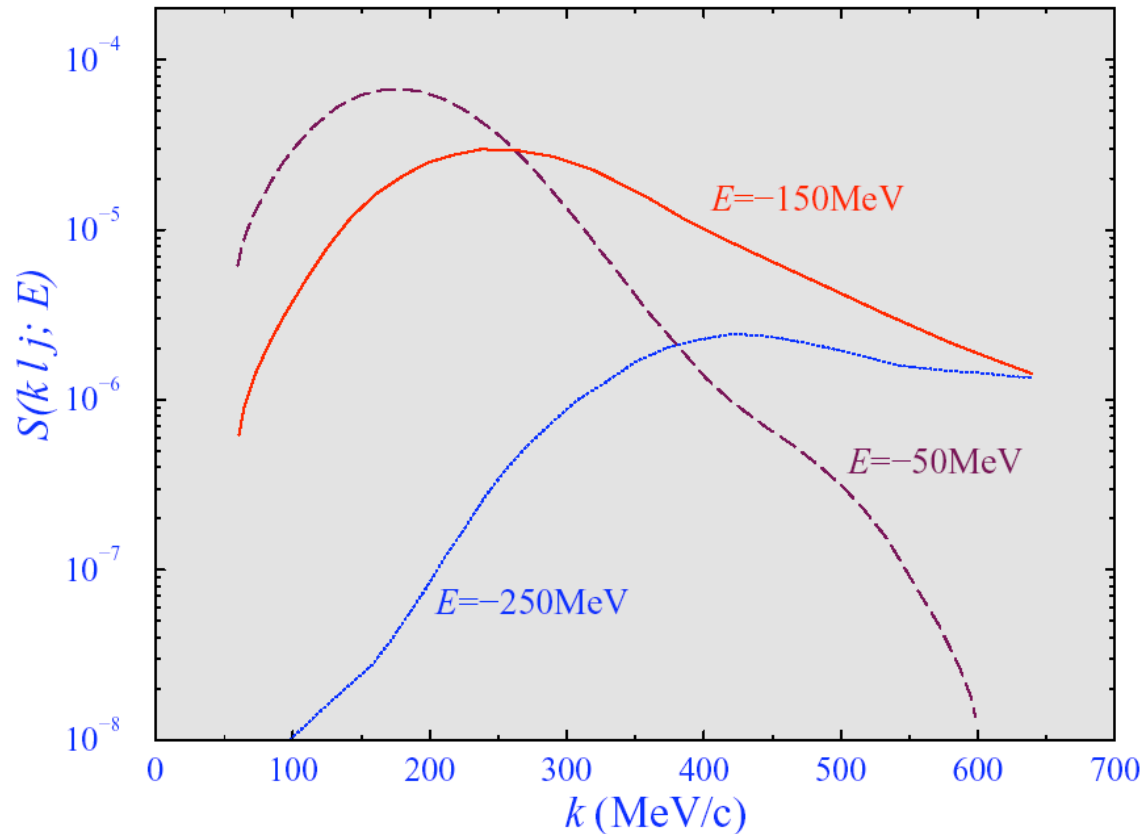
Strength depleted
by 10% due to SRC

Quality of quasihole wave function



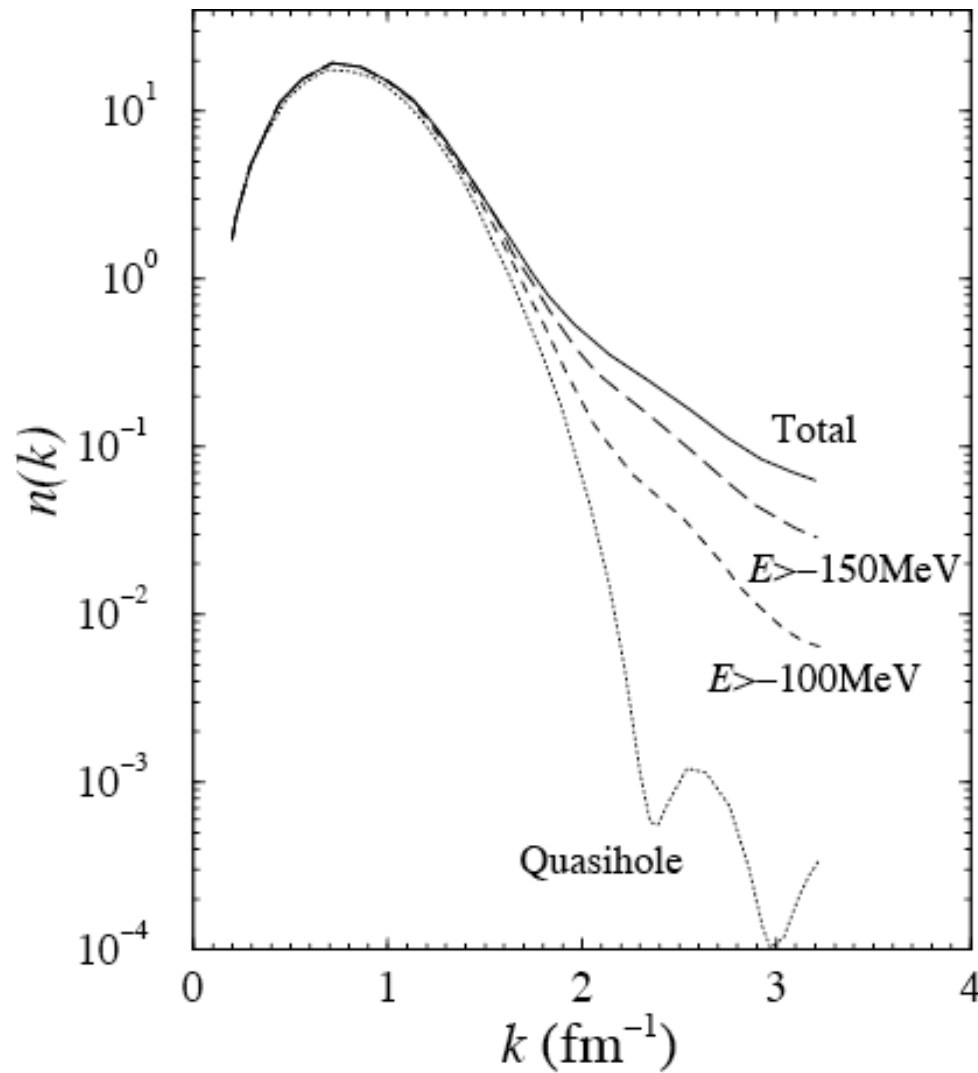
PRC55, 810 (1997)

Prediction of high-momentum components



$p_{1/2}$ spectral function at fixed energies in ^{16}O
Phys. Rev. C49, R17 (1994)

Momentum distribution ^{16}O

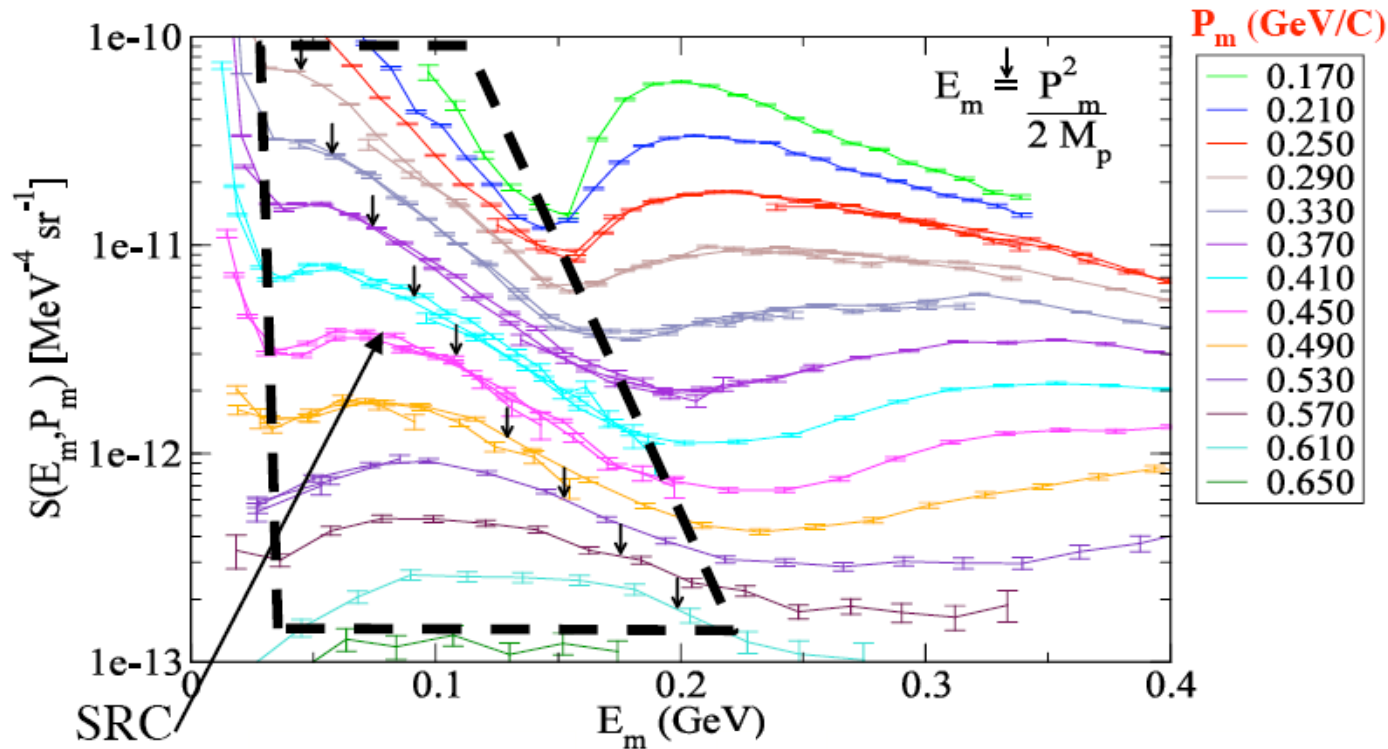


Confirms expectation:

High momentum nucleons
can be found at large
negative energies

What are the rest of the protons doing?

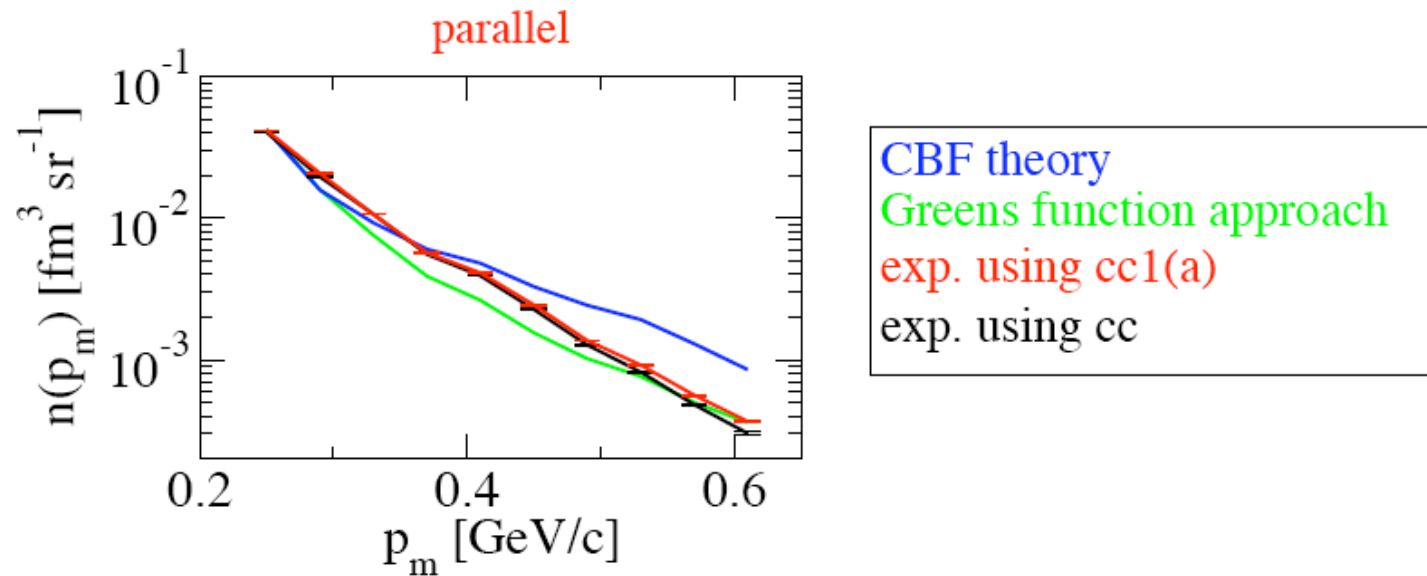
Jlab E97-006 Phys. Rev. Lett. 93, 182501 (2004) D. Rohe et al.



- Location of high-momentum components
- Integrated strength agrees with theoretical prediction Phys. Rev. C49, R17 (1994)
 $\Rightarrow 0.6$ protons for ^{12}C

Integrated strength $\Rightarrow n(k)$

momentum dependence



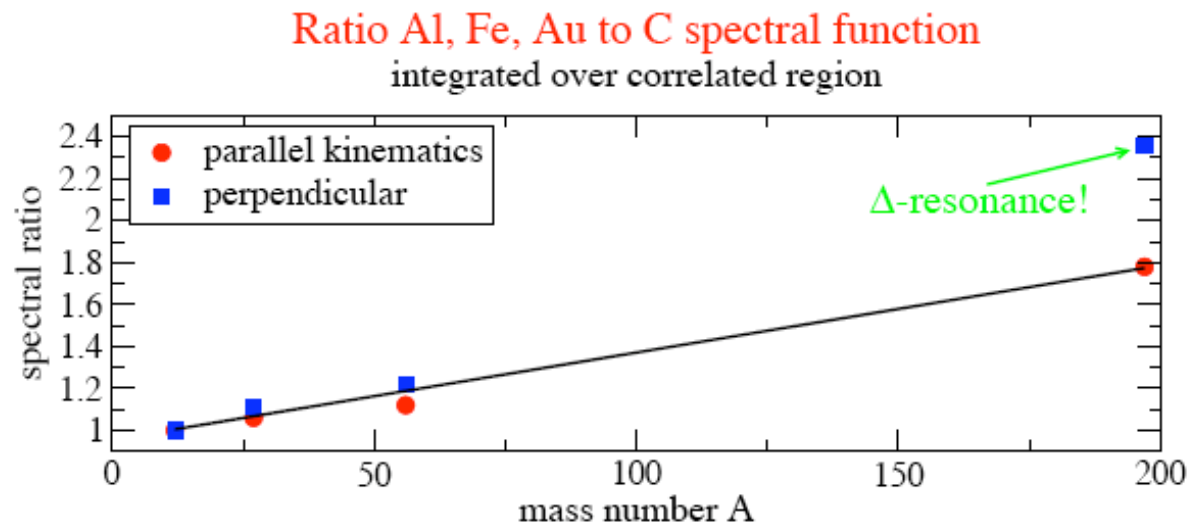
\rightarrow theory and experiment \pm agree

From: Sick, ECT* workshop July 2007

Asymmetry dependence?

ratio to C of correlated strength

find ratio ~ 1 as expected



enhancement for Au

not yet understood

consequence of n-p correlations as $N > Z$??, rescattering ??

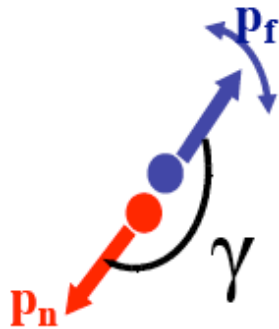
would like to get $S(k, E)$ for $N \neq Z$

Data from BNL

Brookhaven EVA Collaboration Result

A. Tang *et al.*, Phys. Rev. Lett. **90** (2003) 042301.

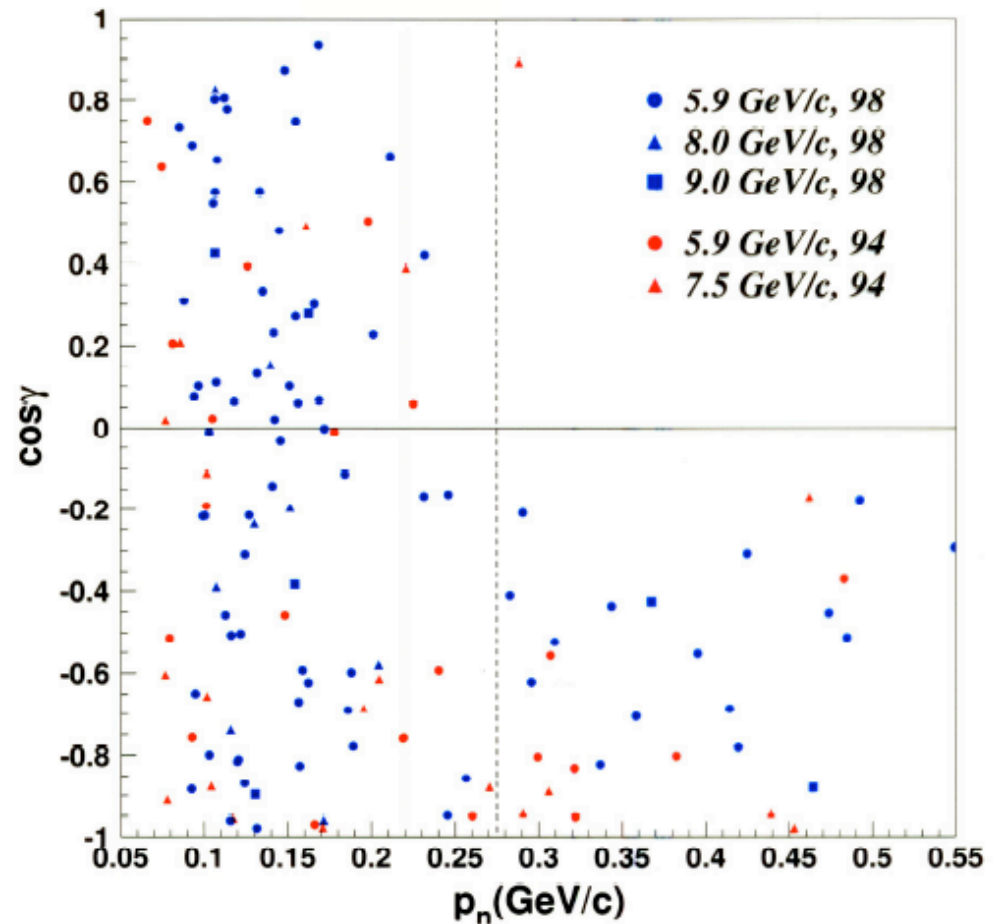
$^{12}\text{C}(p,2p+n)$ Reaction



$$\mathbf{p}_f = \mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_0$$

\mathbf{p}_0 = incident proton

\mathbf{p}_1 and \mathbf{p}_2 are detected

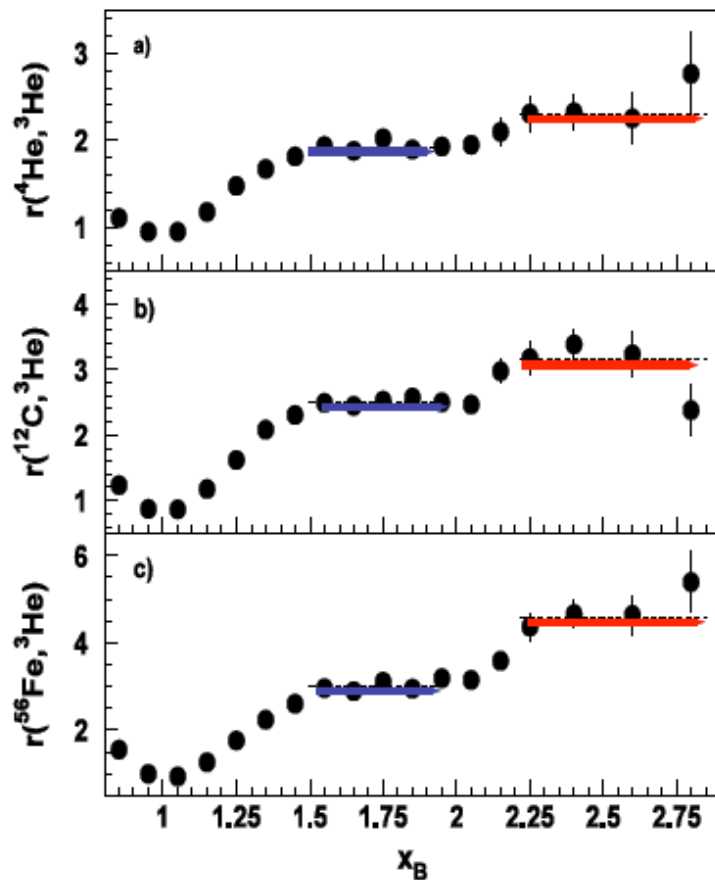


Green's functions
III24

Electron scattering results JLab

New CLAS A(e,e') Result

K. Sh. Egiyan et al., *Phys. Rev. Lett.* 96, 082501 (2006)



The probabilities for 3-nucleon SRC are smaller by one order of magnitude relative to the 2N SRC.

The observed “scaling” means that the electrons probe the high-momentum nucleons in the 3-nucleon phase, and the scaling factors determine the per-nucleon probability of the 3N-SRC phase in nuclei with $A > 3$ relative to ^3He .

Less than 1% of total.

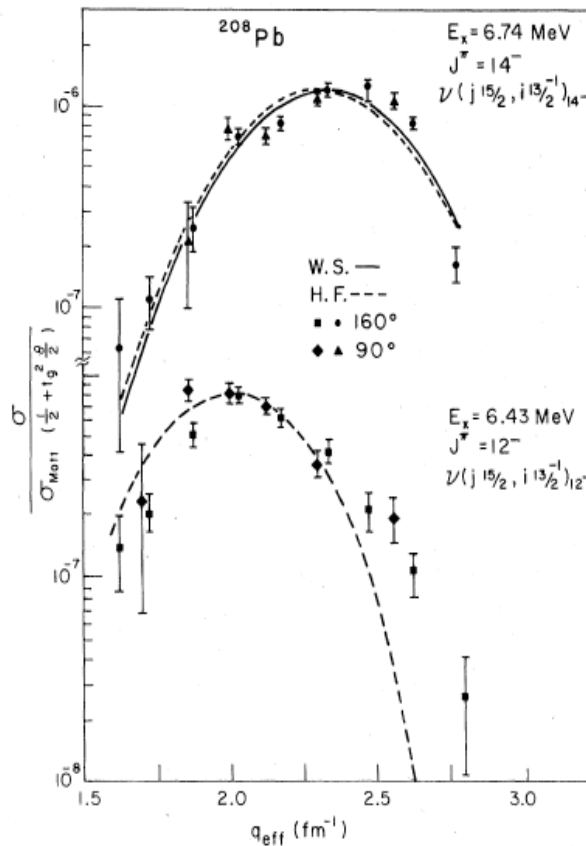
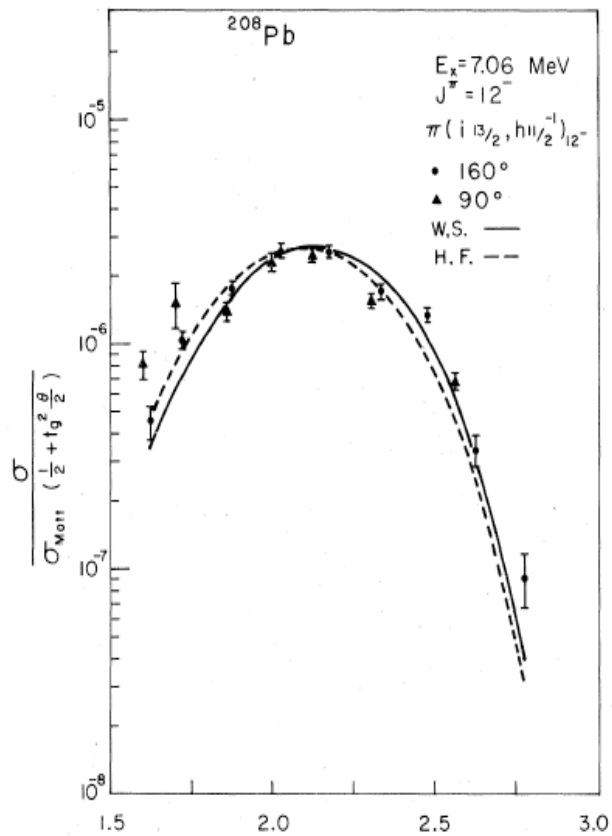
Combination of data & analysis

From the (e,e'), (e,e'p), and (e,e'pN) Results

- 80 +/- 5% single particles moving in an average potential
 - 60 – 70% independent single particle in a shell model potential
 - 10 – 20% shell model long range correlations
- 20 +/- 5% two-nucleon short range correlations *E. Piatsetzky et al., PRL97, 162504 (2006)*
 - from (e,e'pp) approx. 1% pp SRC
 - exact ratio for (e,e'pn) / (e,e'pp) not final, but preliminary result gives approx. 18% pn SRC
 - combining preliminary results we can deduce approx. 1% nn SRC
- less than 1% multi-nucleon correlations

From D.W. Higinbotham, JLab

Consequences of correlated sp strength



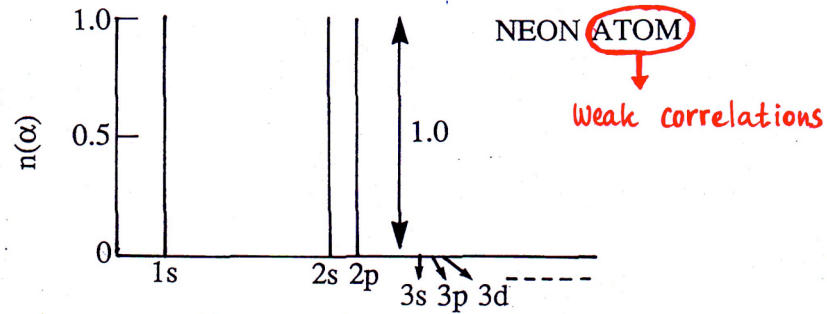
M12 and M14 transitions in (e, e') only 50% of ph estimate

$$\Rightarrow Z_h * Z_p$$

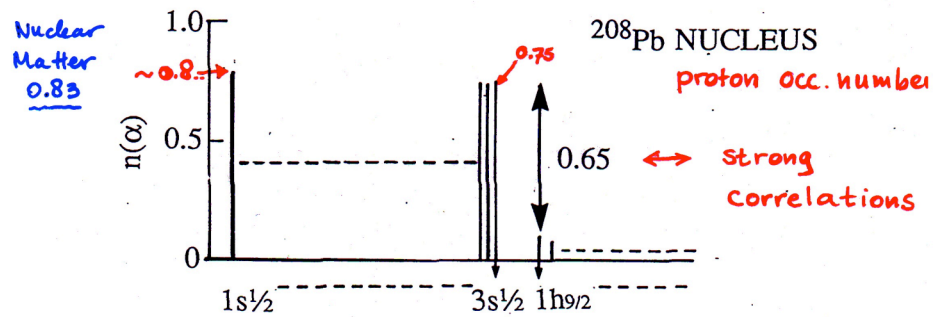
Data:

PRC20, 497(1979)

Occupation Numbers

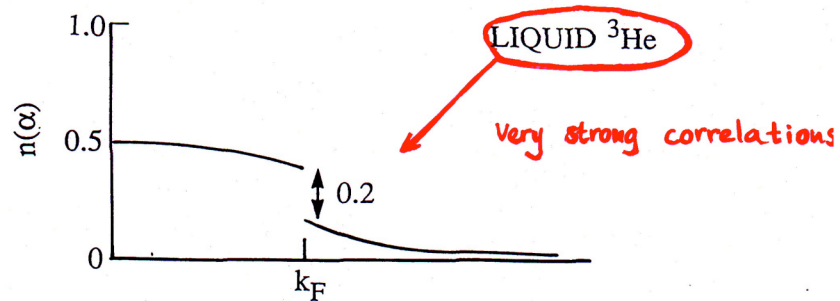


Charge density in ^{208}Pb interior closely linked to nuclear saturation density



In turn, this charge density is predominantly determined by

SHORT-RANGE CORRELATIONS



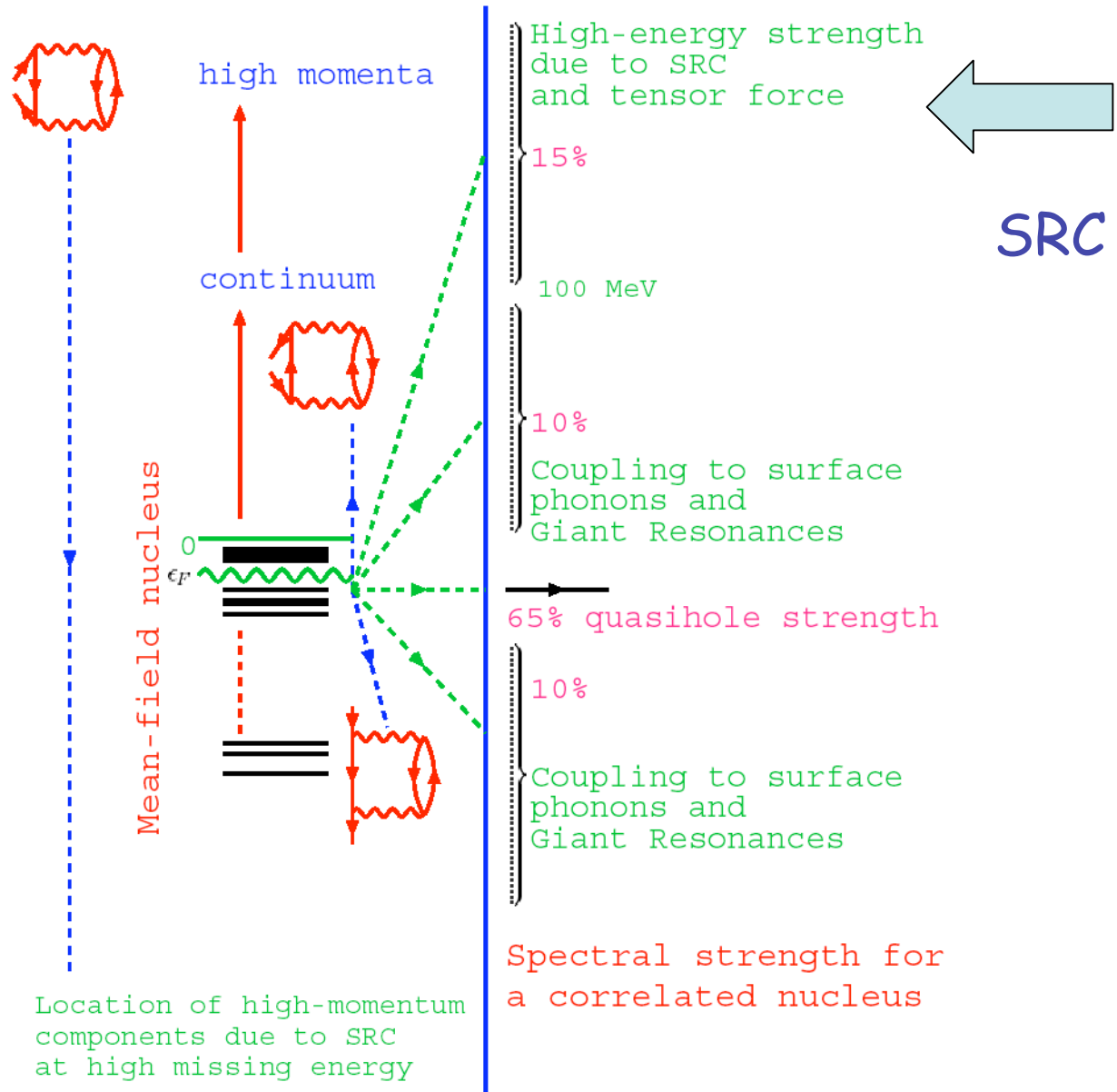
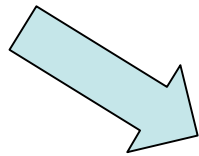
We now essentially know what all the protons are doing in the ground state of a “closed-shell” nucleus !!!

- Unique for a **correlated** many-body system
- Information available for electrons in atoms (Hartree-Fock)
- **Not** for electrons in solids
- **Not** for atoms in quantum liquids
- **Not** for quarks in nucleons

⇒ **Demonstrates the value of the study of the nucleus for its intrinsic interest as a quantum many-body problem!**

Location of single-particle strength in nuclei

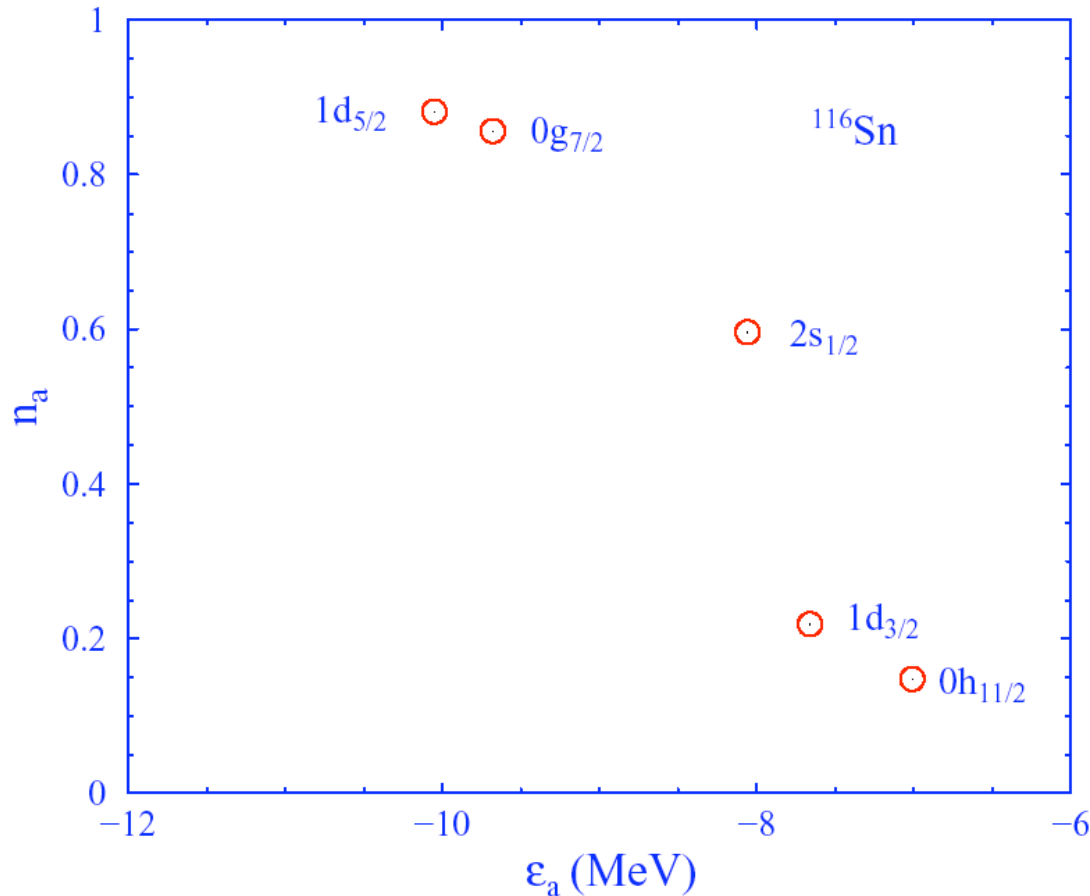
SRC



What about open-shell systems?

Semi-magic nuclei
Green's function calculation

SRC the same
GRs similar



only difference
near ϵ_F

removal &
addition
probabilities
similar size
for $2s_{1/2}$!!

\Rightarrow pairing

Deformation?

$^{142}\text{Nd}(e,e'p)$ $Z=60$; $N=82$ compare with $^{146}\text{Nd}(e,e'p)$ $Z=60$; $N=86$
 Nucl. Phys. **A560**, 811 (1993)

E_x ^{141}Pr

0.000

0.145

1.118

1.298

J^π

$5/2^+$

$7/2^+$

$11/2^-$

$1/2^+$

S_{exp}

0.23

0.39

0.05

0.03

E_x ^{145}Pr

0.000

0.063

0.189

0.348

0.555

J^π

$7/2^+$

$5/2^+$

$5/2^+$

$3/2^+$

$7/2^+$

S_{exp}

0.19

0.17

0.03

0.02

0.03

Wave functions in both
 nuclei are the same!

Systems with N very different from Z ?

SRC still the same (tensor force disappears for n and "increases" for p for $N > Z$)
(see PRC71,014313(2005))

Collectivity of excited states could be reduced

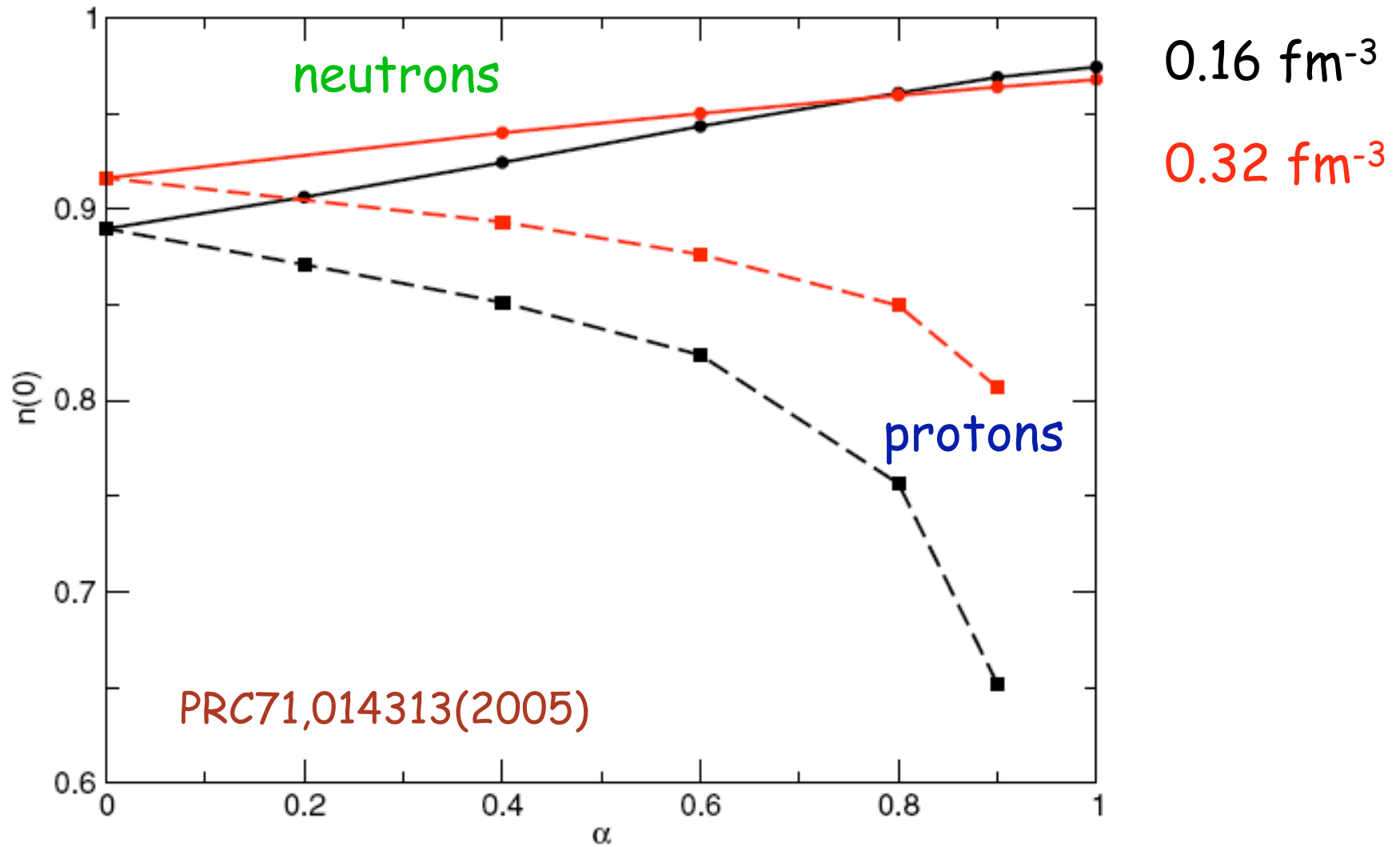
So less fragmentation

and removal of sp strength becomes

more like mean-field (+ SRC+ whatever is left of tensor force for n but perhaps strong effect for p !)

Continuum effects (soft dipoles ...)

SCGF for isospin-polarized nuclear matter



asymmetry